

HÖLDER REGULARITY FOR A KOLMOGOROV EQUATION

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ABSTRACT. We study the interior regularity properties of the solutions to the degenerate parabolic equation,

$$\Delta_x u + b \partial_y u - \partial_t u = f, \quad (x, y, t) \in \mathbb{R}^N \times \mathbb{R} \times \mathbb{R},$$

which arises in mathematical finance and in the theory of diffusion processes.

1. INTRODUCTION

We consider the degenerate parabolic equation

$$(1.1) \quad L_b u \equiv (\Delta_x + b \partial_y - \partial_t) u = f$$

in the variables $(x, y, t) \in \mathbb{R}^N \times \mathbb{R} \times \mathbb{R}$, where Δ_x denotes the Laplacian operator acting in the variables $x = (x_1, \dots, x_N)$. We aim to prove some new Schauder type interior estimates related to the Hölder classes $C_b^{k, \alpha}$ naturally associated to L_b . Our estimates improve the known ones and allow us to study nonlinear equations of the form

$$(1.2) \quad \Delta_x u + b(\cdot, u) \partial_y u - \partial_t u = f(\cdot, u),$$

recently considered in mathematical finance in [1] and [2]. We also obtain regularity results for the following nonlinear convection-diffusion model proposed by Escobedo, Vazquez and Zuazua in [9]:

$$\Delta_x u + \partial_y g(u) - \partial_t u = 0,$$

with particular interest in the case $g(u) = u|u|^{q-1}$ for $q \in]1, \frac{N+2}{N+1}[$.

While we refer to the next section for the precise notation and assumptions on the coefficients b and f , we would like to make some preliminary remarks. One of the main features of operator L_b is the strong degeneracy of its characteristic form due to the lack of diffusion in the y -direction. On the other hand, L_b can be represented in the form,

$$(1.3) \quad \sum_{j=1}^p X_j^2 + X_{p+1},$$

where the first-order differential operators (vector fields) X_j are defined as follows:

$$(1.4) \quad X_j = \partial_{x_j}, \quad j = 1, \dots, p = N, \quad \text{and} \quad X_{N+1} = b \partial_y - \partial_t.$$

A classical result by Hörmander [11] states that if an operator H , in the form (1.3), is such that the vector fields X_j have smooth coefficients and their commutators,

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up to a certain order, span the whole space at every point, then H is hypoelliptic. This means that every weak solution of $Hu = f$, with $f \in C^\infty$, is smooth.

For instance, if $N = 1$ and $b(x, y, t) = x$ in (1.1), then

$$(1.5) \quad L_x = \partial_{xx} + x\partial_y - \partial_t$$

is the linearized prototype of the Kolmogorov operator which, under suitable conditions, describes the probability density of a physical system with two degrees of freedom (cf. [18]). In this case we have

$$X_1 = \partial_x, \quad X_2 = x\partial_y - \partial_t, \quad \text{and} \quad [X_1, X_2] \equiv X_1X_2 - X_2X_1 = \partial_y,$$

so that L_x is a hypoelliptic operator. More generally, the vector fields in (1.4) verify

$$(1.6) \quad [X_j, X_{N+1}] = (\partial_{x_j} b)\partial_y, \quad j = 1, \dots, N;$$

therefore, the assumptions of Hörmander's theorem are satisfied if

$$(1.7) \quad b \in C^\infty \quad \text{and} \quad \nabla_x b \equiv (\partial_{x_1} b, \dots, \partial_{x_N} b) \neq 0.$$

Hörmander's result was the starting point of an extensive study of operators H in the form (1.3) with smooth vector fields. A general theory of the regularity analogous to the classical one has been developed both in Sobolev and Hölder spaces by Folland [10], Rothschild and Stein [17], Nagel and Stein [15], and Beals [3]. We also refer to the more recent papers by Krylov [12] and by Lanconelli, Polidoro and the author [13]. The case of operators in the form

$$\sum_{i,j=1}^p a_{ij} X_i X_j + X_{p+1}$$

with non-regular coefficients a_{ij} has been considered by Xu [19] and Bramanti and Brandolini [4]. Thanks to the known results (cf. [17]), we have the following.

Theorem 1.1 (Rothschild-Stein). *Let u be a classical solution of (1.1) (cf. Definition 4.1) in an open subset Ω of $\mathbb{R}^{N+2}$. If $f \in C_b^{k-2,\alpha}(\Omega)$ and the Hörmander condition (1.7) holds, then $u \in C_b^{k,\alpha}(\Omega)$.*

The regularity assumption on b in Theorem 1.1 can be weakened by assuming at least $b \in C_b^{k+1,\alpha}(\Omega)$. In this case the proof follows the original one with minor changes and we obtain the following:

Theorem 1.2 (Rothschild-Stein). *Let u be a classical solution of (1.1) in Ω with $f \in C_b^{k-2,\alpha}(\Omega)$. If $b \in C_b^{k+1,\alpha}(\Omega)$ and $\nabla_x b \neq 0$ in Ω , then $u \in C_b^{k,\alpha}(\Omega)$.*

In view of the classical Schauder estimates, the previous results do not seem optimal. In particular, we emphasize that they *do not* allow the treatment of the existence and regularity theory of nonlinear equations. As a matter of fact, the further weaker assumption $b \in C_b^{k-2,\alpha}$ is naturally expected. Actually, the techniques used by Rothschild and Stein require the smoothness of the vector fields as an essential hypothesis. On the contrary, here we aim to consider *non-regular* vector fields.

In the recent papers [7], [8] in collaboration with Citti and Polidoro, we considered the nonlinear equation in three variables

$$(1.8) \quad L_u u = \partial_{xx} u + u\partial_y u - \partial_t u = f$$

and we studied the regularity of the solution u by a modification of the classical freezing method. More precisely, we regarded L_b as a local perturbation of a Hörmander's operator on the Heisenberg group. This last operator played the same role as the constant coefficients operators in the classical theory. This technique was introduced by Citti in [5] to study an equation of Levi type.

Aiming to adapt those ideas, we immediately realize that, in dimensions higher than three, the Lie algebra formally associated to L_b is not free. This means that the vector fields X_j do not satisfy as few linear relations as possible (i.e., only those forced by anti-commutativity and the Jacobi identity). As a consequence, the algebra that one might naturally associate to L_b varies from point to point. In order to overcome this problem and to eliminate the inessential relations among the commutators, we add some extra variables and we lift the operator L_b to a higher-dimensional space. We recall that a general version of the so-called "lifting method" for an operator in (1.3) with *smooth* coefficients, is due to Rothschild and Stein [17]. In our case we make the tentative choice to define the following operator in $\mathbb{R}^{2N+1}$:

$$(1.9) \quad L_B = \Delta_x + b\partial_{y_1} + x_2\partial_{y_2} + \dots + x_N\partial_{y_N} - \partial_t$$

where $(x, y, t) = (x_1, \dots, x_N, y_1, \dots, y_N, t)$ denotes the point in $\mathbb{R}^{2N+1}$. In order to apply to (1.9) the freezing techniques cited above, a detailed analysis and careful estimates of the fundamental solutions to the frozen operators are in order. This is done in Section 3 and it is our main proof. Then, we study the regularity properties of L_B and finally, we apply our results to the operator L_b . We prove the following.

Theorem 1.3. *Let u be a classical solution of (1.1) in Ω (cf. Definition 4.1 and Remark 4.3) and assume the Hörmander type condition*

$$(1.10) \quad b \in C^1(\Omega) \quad \text{and} \quad \nabla_x b \neq 0 \quad \text{in } \Omega.$$

If $b, f \in C_b^{k-2, \alpha}(\Omega)$, with $k \geq 2$ and $\alpha \in]0, 1[$, then $u \in C_b^{k, \bar{\alpha}}(\Omega)$ for every $\bar{\alpha} \in]0, \alpha[$.

We remark that Theorem 1.3 and a bootstrap argument give simple conditions for the interior regularity of solutions to a nonlinear equation of the form (1.2). In particular, we refine the results in [7].

Two possible directions for extending Theorem 1.3 come readily to mind. It seems that our technique can be adapted without difficulty to the following, more general, class of ultraparabolic operators in $\mathbb{R}^{N+2}$:

$$\sum_{i,j=1}^N a_{ij} \partial_{x_i x_j} + b \partial_y - \partial_t$$

where (a_{ij}) is a positive definite matrix with Hölder continuous entries. Secondly, assumption (1.10) could be relaxed by a "higher step" condition, that is, by requiring that higher-order commutators of the vector fields X_j span $\mathbb{R}^{N+2}$. In this case, it seems that the proof would be essentially analogous, even if it could become considerably knotty.

This paper is organized as follows. In Section 2 we set the notation and we collect some tools for the analysis on nilpotent Lie groups. In Section 3 we provide some estimates of the fundamental solutions of the frozen operators. Section 4 is devoted to the proof of Theorem 1.3.

2. HÖLDER CLASSES AND CONTROL DISTANCES

In this section we present some preliminary material and we define the lifted and frozen operators related to L_b . We begin by defining the Hölder classes related to the vector fields in (1.4). For the reader's convenience, we also give the following standard

Definition 2.1. Let D be a locally Lipschitz continuous vector field on Ω and d a (positive) formal degree associated to D . We call u Hölder continuous with exponent α , $\alpha \in]0, d[$, in Ω w.r.t. D and we write $u \in C_D^\alpha(\Omega)$ if, for every compact subset E of Ω , there exists a constant C such that

$$|u(\exp(\delta D)(z)) - u(z)| \leq C|\delta|^{\frac{\alpha}{d}},$$

for every $z \in E$ and suitably small δ . We refer, for instance, to [16] for the definition and properties of exponential mappings induced by vector fields. We say that u is Lie derivable w.r.t. D in $z \in \Omega$ if the following limit exists:

$$Du(z) = \lim_{\delta \rightarrow 0} \frac{u(\exp(\delta D)(z)) - u(z)}{\delta}.$$

Definition 2.2. Let $\alpha \in]0, 1[$. We set the formal degree of $X_1, \dots, X_N$ in (1.4) equal to one and the formal degree of X_{N+1} equal to two. We define

$$C_b^\alpha(\Omega) = \bigcap_{j=1}^{N+1} C_{X_j}^\alpha(\Omega).$$

We say that $u \in C_b^{1,\alpha}(\Omega)$ if

$$X_j u \in C_b^\alpha(\Omega), \quad j = 1, \dots, N, \quad \text{and} \quad u \in C_{X_{N+1}}^{1+\alpha}(\Omega).$$

Finally, if $k \in \mathbb{N}$, $k \geq 2$, we define by recurrence the class $C_b^{k,\alpha}(\Omega)$ as follows: assuming that $b \in C_b^{k-2,\alpha}(\Omega)$, we say that $u \in C_b^{k,\alpha}(\Omega)$ if

$$X_j u \in C_b^{k-1,\alpha}(\Omega), \quad j = 1, \dots, N, \quad \text{and} \quad X_{N+1} u \in C_b^{k-2,\alpha}(\Omega).$$

For greater convenience, when in the sequel we consider the class $u \in C_b^{k,\alpha}(\Omega)$ with $k \geq 2$, we always implicitly assume that $b \in C_b^{k-2,\alpha}(\Omega)$.

No regularity in the y -direction is seemingly assumed in the definition of $C_b^{k,\alpha}(\Omega)$. On the other hand, keeping in mind that $[X_j, X_{N+1}] = (\partial_{x_j} b) \partial_y$ and Hörmander's condition, it is natural to set the formal degree of the vector field ∂_y equal to three and it is possible to prove, by a standard argument based on the Campbell-Hausdorff formula, the following

Lemma 2.3. *If $k = 0, 1, 2$ and $u \in C_b^{k,\alpha}(\Omega)$, then $u \in C_{\partial_y}^{k+\alpha}(\Omega)$. If $k \geq 3$ and $u \in C_b^{k,\alpha}(\Omega)$, then $\partial_y u \in C_b^{k-3,\alpha}(\Omega)$.*

By the previous lemma, we have the following inclusion of the space $C_b^{k,\alpha}$ in the space of Hölder continuous functions in the classical sense:

$$C_b^{3k,\alpha}(\Omega) \subseteq C^{k, \frac{\alpha}{3}}(\Omega).$$

We now lift the original vector fields in (1.4) to $\mathbb{R}^{2N+1}$ in such a way that they become free. Since we aim to prove a local result, it is not restrictive to suppose that Ω is suitably small. Then, without loss of generality, by (1.10), we may assume that

$\partial_{x_1} b \neq 0$ in Ω . In the sequel we denote by $z = (x, y, t) = (x_1, \dots, x_N, y_1, \dots, y_N, t)$ and $\zeta = (\xi, \eta, \tau)$ the points in $\mathbb{R}^{2N+1}$ and we set $\Omega_0 = \Omega \times \mathbb{R}^{N-1}$.

We define the lifted vector fields on Ω_0 as

$$(2.1) \quad D_j = \partial_{x_j}, \quad j = 1, \dots, N, \quad D_{N+1} = \langle B, \nabla_y \rangle - \partial_t,$$

where

$$B(x, y_1, t) = (b(x, y_1, t), x_2, \dots, x_N), \quad \text{and} \quad \nabla_y = (\partial_{y_1}, \dots, \partial_{y_N}).$$

Thus, the operator L_B in (1.9) can be expressed in the form

$$L_B = \sum_{j=1}^N D_j^2 + D_{N+1}.$$

Since $\partial_{x_1} b \neq 0$ in Ω_0 , the commutators

$$D_{N+2} \equiv [D_1, D_{N+1}] = (\partial_{x_1} b) \partial_{y_1},$$

$$D_{N+1+j} \equiv [D_j, D_{N+1}] = (\partial_{x_j} b) \partial_{y_j} + \partial_{y_j}, \quad 2 \leq j \leq N,$$

are linearly independent and the system $(D_j)_{1 \leq j \leq 2N+1}$ forms a basis of $\mathbb{R}^{2N+1}$ at every point of Ω_0 . Analogously to Definition 2.2, we give the notion of Hölder continuity related to (D_j) .

Definition 2.4. Let $\alpha \in]0, 1[$. We set the formal degree of $D_1, \dots, D_N$ equal to one and the formal degree of D_{N+1} equal to two. We define

$$C_B^\alpha(\Omega_0) = \bigcap_{j=1}^{N+1} C_{D_j}^\alpha(\Omega_0).$$

We say that $u \in C_B^{1,\alpha}(\Omega_0)$ if

$$D_j u \in C_B^\alpha(\Omega_0), \quad j = 1, \dots, N, \quad \text{and} \quad u \in C_{D_{N+1}}^{1+\alpha}(\Omega_0).$$

Finally, if $k \in \mathbb{N}$, $k \geq 2$, we define by recurrence the class $C_B^{k,\alpha}(\Omega_0)$ as follows: assuming that $b \in C_b^{k-2,\alpha}(\Omega)$, we say that $u \in C_B^{k,\alpha}(\Omega_0)$ if

$$D_j u \in C_B^{k-1,\alpha}(\Omega_0), \quad j = 1, \dots, N, \quad \text{and} \quad D_{N+1} u \in C_B^{k-2,\alpha}(\Omega_0).$$

Remark 2.5. Given a function $w = w(x, y_1, t)$ on Ω , we denote again by w its extension to $\Omega_0 = \Omega \times \mathbb{R}^{N-1}$, i.e., the function defined by $w(x, y_1, \dots, y_N, t) = w(x, y_1, t)$. Hence, it is clear that a solution u to (1.1) in Ω is also a solution to $L_B u = f$ in Ω_0 . Moreover, $u \in C_b^{k,\alpha}(\Omega)$ if and only if $u \in C_B^{k,\alpha}(\Omega_0)$.

We next construct a nilpotent Hörmander operator locally approximating L_B and we introduce some distances naturally associated to the vector fields D_j in (2.1). More details about distances defined by vector fields can be found in [10] and [16].

For fixed $\bar{z} \in \Omega_0$, we define the frozen vector fields

$$(2.2) \quad D_j^{\bar{z}} = \partial_{x_j}, \quad j = 1, \dots, N,$$

and

$$D_{N+1}^{\bar{z}} = (b(\bar{z}) + \langle \nabla_x b(\bar{z}), (x - \bar{x}) \rangle) \partial_{y_1} + \sum_{j=2}^N x_j \partial_{y_j} - \partial_t.$$

Since the commutators of $D_j^{\bar{z}}$ and $D_{N+1}^{\bar{z}}$ are given by

$$(2.3) \quad \begin{aligned} D_{N+2}^{\bar{z}} &\equiv [D_1^{\bar{z}}, D_{N+1}^{\bar{z}}] = \partial_{x_1} b(\bar{z}) \partial_{y_1}, \\ D_{N+1+j}^{\bar{z}} &\equiv [D_j^{\bar{z}}, D_{N+1}^{\bar{z}}] = \partial_{x_j} b(\bar{z}) \partial_{y_1} + \partial_{y_j}, \quad 2 \leq j \leq N, \end{aligned}$$

and $\partial_{x_1} b(\bar{z}) \neq 0$ by assumption, the Hörmander condition is verified and the operator

$$(2.4) \quad L_{\bar{z}} = \sum_{j=1}^N (D_j^{\bar{z}})^2 + D_{N+1}^{\bar{z}}$$

is hypoelliptic. We call

$$\nabla_{\bar{z}} \equiv (D_1^{\bar{z}}, \dots, D_{2N+1}^{\bar{z}}),$$

the intrinsic gradient related to the system of vector fields defined in (2.2) and (2.3). For fixed $z \in \mathbb{R}^{2N+1}$, we consider the exponential map

$$E_z^{\bar{z}}(\theta) = \exp(\langle \theta, \nabla_{\bar{z}} \rangle)(z), \quad \theta \in \mathbb{R}^{2N+1}.$$

It is well known that the map $E_z^{\bar{z}}$ is a global diffeomorphism. Its inverse function $\theta_z^{\bar{z}}$ is usually called the canonical change of coordinates and it has the explicit expression

$$(2.5) \quad \theta_z^{\bar{z}}(\zeta) = (E_z^{\bar{z}})^{-1}(\zeta) = (\theta_1, \dots, \theta_{2N+1}),$$

where

$$(2.6) \quad \begin{aligned} \theta_j &= \xi_j - x_j, \quad 1 \leq j \leq N, \\ \theta_{N+1} &= -(\tau - t), \\ \theta_{N+2} &= \frac{1}{\partial_{x_1} b(\bar{z})} \left[\eta_1 - y_1 + (\tau - t) \left(b(\bar{z}) + \frac{\partial_{x_1} b(\bar{z})}{2} (\xi_1 + x_1 - 2\bar{x}_1) \right) \right. \\ &\quad \left. - \sum_{j=2}^N \partial_{x_j} b(\bar{z}) (\eta_j - y_j + \bar{x}_j (\tau - t)) \right], \\ \theta_{N+1+j} &= \eta_j - y_j + \frac{(\tau - t)(\xi_j + x_j)}{2}, \quad 2 \leq j \leq N. \end{aligned}$$

Through the canonical change of coordinates, the vector fields $D_j^z, D_j^{\bar{z}}$ corresponding to different points $z, \bar{z} \in \Omega_0$, coincide. More precisely, if we set

$$(2.7) \quad \begin{aligned} D_j^H &= \partial_{\theta_j} - \frac{\theta_{N+1}}{2} \partial_{\theta_{N+1+j}}, \quad 1 \leq j \leq N, \\ D_{N+1}^H &= \partial_{\theta_{N+1}} + \frac{1}{2} \sum_{j=1}^N \theta_j \partial_{\theta_{N+1+j}}, \\ D_{N+1+j}^H &\equiv [D_j^H, D_{N+1}^H] = \partial_{\theta_{N+1+j}}, \quad 1 \leq j \leq N, \end{aligned}$$

then, for any smooth function φ and $\bar{z} \in \Omega_0$, it follows that

$$D_j^{\bar{z}}(\varphi \circ \theta_z^{\bar{z}}) = (D_j^H \varphi) \circ \theta_z^{\bar{z}}, \quad 1 \leq j \leq 2N+1.$$

The vector fields in (2.7) generate a free Lie algebra which is isomorphic to the Heisenberg one. Indeed, the vector fields in (2.7) induce a composition law in $\mathbb{R}^{2N+1}$ formally defined by the Campbell-Hausdorff formula, or explicitly

$$(\theta \oplus \bar{\theta})_j = \begin{cases} \theta_j + \bar{\theta}_j, & \text{if } 1 \leq j \leq N+1, \\ \theta_j + \bar{\theta}_j + \frac{1}{2} (\theta_{j-N-1} \bar{\theta}_{N+1} - \bar{\theta}_{j-N-1} \theta_{N+1}), & \text{if } N+2 \leq j \leq 2N+1, \end{cases}$$

and the dilations group

$$(2.8) \quad \delta_\lambda(\theta) = (\lambda\theta_1, \dots, \lambda\theta_N, \lambda^2\theta_{N+1}, \lambda^3\theta_{N+2}, \dots, \lambda^3\theta_{2N+1}), \quad \lambda > 0.$$

The space $\mathbb{R}^{2N+1}$ endowed with the law $\oplus$ and the dilations δ_λ is a homogeneous Lie group. The associated Lie algebra of the $\oplus$ -left-invariant vector fields is the one generated by $D_1^H, \dots, D_{N+1}^H$. We also remark that $D_1^H, \dots, D_N^H$ are δ_λ -homogeneous of degree one and D_{N+1}^H is δ_λ -homogeneous of degree two. Therefore, the Hörmander operator,

$$(2.9) \quad L_H = \sum_{j=1}^N (D_j^H)^2 + D_{N+1}^H,$$

has a fundamental solution Γ_H which is invariant with respect to the left $\oplus$ -translations and it is homogeneous of degree $-Q + 2$ where $Q = 4N + 2$ is the homogeneous dimension of $(\mathbb{R}^{N+1}, \oplus)$. An explicit expression of Γ_H is known (see, e.g., [14]); however, here we only use its qualitative properties. A norm homogeneous w.r.t. the dilations in (2.8) is given by

$$(2.10) \quad \|\theta\|_H = \sum_{j=1}^N \left(|\theta_j| + |\theta_{N+1+j}|^{\frac{1}{3}} \right) + |\theta_{N+1}|^{\frac{1}{2}},$$

and the associated control distance is defined by

$$d_H(\bar{\theta}, \theta) = \|\theta^{-1} \oplus \bar{\theta}\|_H.$$

The following Lie product on $\mathbb{R}^{2N+1}$ is naturally induced:

$$(2.11) \quad \begin{aligned} z \circ \zeta &= E_{\bar{z}}^{\bar{z}}(\theta_{\bar{z}}^{\bar{z}}(z) \oplus \theta_{\bar{z}}^{\bar{z}}(\zeta)) \\ &= (x + \xi - \bar{x}, y + \eta - \bar{y} - (\tau - \bar{t})\mathcal{J}_x B(\bar{z})(x - \bar{x}), t + \tau - \bar{t}), \end{aligned}$$

where $\mathcal{J}_x B$ denotes the Jacobian matrix of B w.r.t. the variable x , i.e., the diagonal matrix $\text{diag}(\partial_{x_1} b, 1, \dots, 1)$. Correspondingly, we have the dilations group

$$\delta_\lambda^{\bar{z}}(z) = E_{\bar{z}}^{\bar{z}}(\delta_\lambda \theta_{\bar{z}}^{\bar{z}}(z)), \quad \lambda > 0,$$

and the associated control distance

$$(2.12) \quad d_{\bar{z}}(z, \zeta) = \|\theta_{\bar{z}}^{\bar{z}}(\zeta)\|_H.$$

Lemma 2.6. *There exists a constant c_0 , only dependent on Ω_0 , such that*

$$(2.13) \quad c_0^{-1} d_z(z, \zeta) \leq \tilde{d}_z(z, \zeta) \leq c_0 d_z(z, \zeta),$$

for every $z, \zeta \in \Omega_0$, where

$$\begin{aligned} \tilde{d}_z(z, \zeta) &= |x - \xi| + |t - \tau|^{\frac{1}{2}} + |y_1 - \eta_1 + (t - \tau)b(z)|^{\frac{1}{3}} \\ &\quad + \sum_{j=2}^N \left| y_j - \eta_j + \frac{(t - \tau)(\xi_j + x_j)}{2} \right|^{\frac{1}{3}}. \end{aligned}$$

Proof. By (2.6) and by denoting $\theta_z^z(\zeta) = (\theta_1, \dots, \theta_{2N+1})$, we have

$$\begin{aligned} \theta_{N+2} &= \frac{1}{\partial_{x_1} b(z)} \left[\eta_1 + y_1 + b(z)(\tau - t) - \sum_{j=2}^N \partial_{x_j} b(z) \left(\eta_j - y_j + \frac{(\tau - t)(\xi_j + x_j)}{2} \right) \right. \\ &\quad \left. + \frac{1}{2} \sum_{j=2}^N \partial_{x_j} b(z)(\tau - t)(\xi_j - x_j) \right] + \frac{(\tau - t)(\xi_1 - x_1)}{2} \\ &= \frac{1}{\partial_{x_1} b(z)} \left[\eta_1 + y_1 + b(z)(\tau - t) - \sum_{j=2}^N \partial_{x_j} b(z) \left(\theta_{N+1+j} + \frac{1}{2} \theta_j \theta_{N+1} \right) \right] \\ &\quad - \frac{1}{2} \theta_1 \theta_{N+1}. \end{aligned}$$

Hence (2.13) is an immediate consequence of the elementary inequality

$$(2.14) \quad (ab)^{\frac{1}{3}} \leq \frac{a}{3} + \frac{2\sqrt{b}}{3}, \quad \forall a, b > 0.$$

□

We stress that the distances d_z, d_ζ corresponding to different points $z, \zeta \in \Omega_0$ are not, in general, equivalent. Nevertheless, using Lemma 2.6, it is straightforward to prove the following

Lemma 2.7. *There exists a constant c_0 , only dependent on Ω_0 , such that*

$$(2.15) \quad c_0^{-1} d_z(z, \bar{z}) \leq d_{\bar{z}}(\bar{z}, z) \leq c_0 d_z(z, \bar{z})$$

and

$$(2.16) \quad d_z(z, \zeta) \leq c_0 (d_z(z, \bar{z}) + d_{\bar{z}}(\bar{z}, \zeta))$$

for every $z, \bar{z}, \zeta \in \Omega_0$.

It is remarkable that the Hölder continuity property related to the vector fields D_j can be expressed in terms of the control distances associated to the frozen vector fields. Indeed, we have

Lemma 2.8. *Let g be a function on Ω_0 and $\alpha \in]0, 1[$. Suppose that, for every compact subset E of Ω_0 , there exists a constant C such that*

$$(2.17) \quad |g(z) - g(\bar{z})| \leq C d_{\bar{z}}(\bar{z}, z)^\alpha, \quad \forall z, \bar{z} \in E.$$

Then $g \in C_B^\alpha(\Omega_0)$.

Proof. We have to prove that

$$(2.18) \quad \begin{aligned} |u(\exp(\delta D_j)(z)) - u(z)| &\leq c\delta^\alpha, \quad 1 \leq j \leq N, \\ |u(\exp(\delta D_{N+1})(z)) - u(z)| &\leq c\delta^{\frac{\alpha}{2}}. \end{aligned}$$

The first inequality is obvious since $d_z(z, \exp(\delta D_j)(z)) = \delta$, for $1 \leq j \leq N$. With regard to the second inequality in (2.18), by assumption (2.17), it suffices to verify that

$$(2.19) \quad d_z(z, \exp(\delta D_{N+1})(z)) \leq c\delta^{\frac{1}{2}}.$$

Denoting $\gamma(\delta) = \exp(\delta D_{N+1})(z)$, we have

$$\begin{aligned}\gamma(\delta) &= z + \int_0^\delta D_{N+1}(\gamma(s))ds \\ &= \left(x, y_1 + \int_0^\delta b(\gamma(s))ds, y_2 + \delta x_2, \dots, y_N + \delta x_N, t - \delta \right).\end{aligned}$$

Hence, by Lemma 2.6, we get

$$\begin{aligned}c_0^{-1}d_z(z, \gamma(\delta)) &\leq \tilde{d}_z(z, \gamma(\delta)) = |\delta|^{\frac{1}{2}} + \left| \int_0^\delta (b(\gamma(s)) - b(z))ds \right|^{\frac{1}{3}} \\ &\leq |\delta|^{\frac{1}{2}} + \left| \delta \sup_{s \in [0, \delta]} |b(\gamma(s)) - b(z)| \right|^{\frac{1}{3}}\end{aligned}$$

(since $b \in C^1(\Omega)$ and $|z - \zeta| \leq d_z(z, \zeta)$ for every z, ζ in a suitably compact subset E_0 of Ω)

$$\leq |\delta|^{\frac{1}{2}} + \left| \delta \left(\sup_{E_0} |\nabla b| \right) d_z(z, \gamma(\delta)) \right|^{\frac{1}{3}}$$

(by (2.14) for $\varepsilon > 0$)

$$\leq |\delta|^{\frac{1}{2}} \left(1 + \frac{2}{3} \left(\varepsilon \sup_{E_0} |\nabla b| \right)^{\frac{1}{2}} \right) + \frac{d_z(z, \gamma(\delta))}{3\varepsilon},$$

which yields (2.19) if ε is suitably large. $\square$

The control distances previously introduced also give an estimate of the error in the intrinsic Taylor expansion of a function $u \in C_B^{k, \alpha}$. To be more precise, as in [7], Theorem 2.16, the following result can be proved.

Proposition 2.9. *Let $\bar{z} \in \Omega_0$ and $u \in C_B^{k, \alpha}(\Omega_0)$. There exists a unique polynomial function $P_{\bar{z}}^k u$ which is a sum of terms $\delta_{\lambda}^{\bar{z}}$ -homogeneous of degree less than or equal to k and verifies*

$$(2.20) \quad u(z) = P_{\bar{z}}^k u(z) + O(d_{\bar{z}}(\bar{z}, z)^{k+\alpha}), \quad \text{as } z \longrightarrow \bar{z}.$$

For instance, in the case $k = 0, 1$, we have

$$P_{\bar{z}}^0 u(z) = u(\bar{z}) \quad \text{and} \quad P_{\bar{z}}^1 u(z) = u(\bar{z}) + \langle \nabla_x u(\bar{z}), x - \bar{x} \rangle.$$

Hence, the frozen vector fields defined in (2.2) are obtained by considering the first-order (intrinsic) Taylor expansion of the coefficients of the original vector fields in (2.1). In particular, we have

$$D_{N+1} - D_{N+1}^{\bar{z}} = (b - P_{\bar{z}}^1 b) \partial_{y_1}.$$

We end this section by stating a technical lemma, which will be used in the proof of Theorem 1.3. For fixed $z, \bar{z} \in \Omega_0$ and a constant $M \geq 1$, we define the set

$$(2.21) \quad \Omega_M(\bar{z}, z) = \{\zeta \in \Omega_0 \mid M d_{\bar{z}}(\bar{z}, z) \leq d_{\bar{z}}(\bar{z}, \zeta)\}.$$

We remark that we can choose M sufficiently large so that

$$(2.22) \quad (c_M)^{-1} d_z(z, \zeta) \leq d_{\bar{z}}(\bar{z}, \zeta) \leq c_M d_z(z, \zeta), \quad \forall \zeta \in \Omega_M(\bar{z}, z),$$

for some constant c_M . Indeed, by Lemma 2.7, we have

$$d_z(z, \zeta) \leq c_0 (d_z(z, \bar{z}) + d_{\bar{z}}(\bar{z}, \zeta)) \leq c_0 (c_0 d_{\bar{z}}(\bar{z}, z) + d_{\bar{z}}(\bar{z}, \zeta)) \leq c_0 \left(\frac{c_0}{M} + 1 \right) d_{\bar{z}}(\bar{z}, \zeta),$$

and, on the other hand,

$$d_{\bar{z}}(\bar{z}, \zeta) \leq c_0 (d_{\bar{z}}(\bar{z}, z) + d_z(z, \zeta)) \leq c_0 \left(\frac{d_{\bar{z}}(\bar{z}, \zeta)}{M} + d_z(z, \zeta) \right),$$

so that

$$\left(\frac{1}{c_0} - \frac{1}{M} \right) d_{\bar{z}}(\bar{z}, \zeta) \leq d_z(z, \zeta).$$

Lemma 2.10. *Let $u \in C_B^{k,\alpha}(\Omega_0)$ and $k \leq 4$. For every compact subset E of Ω_0 , there exists a constant $c = c(E)$ such that*

$$(2.23) \quad |P_z^k u(\zeta) - P_{\bar{z}}^k u(\zeta)| \leq c d_{\bar{z}}(\bar{z}, z)^\alpha d_{\bar{z}}(\bar{z}, \zeta)^k,$$

$$(2.24) \quad |(P_z^k u(\zeta) - P_z^1 u(\zeta)) - (P_{\bar{z}}^k u(\zeta) - P_{\bar{z}}^1 u(\zeta))| \leq c d_{\bar{z}}(\bar{z}, z)^\alpha d_{\bar{z}}(\bar{z}, \zeta)^k,$$

for $z, \bar{z} \in E$ and $\zeta \in \Omega_M(\bar{z}, z)$.

Proof. The proof is a direct and tiresome computation. We only show (2.23) for $k = 1$. We have

$$|P_z^1 u(\zeta) - P_{\bar{z}}^1 u(\zeta)| = |u(z) - P_{\bar{z}}^1 u(z) + \langle \xi - x, \nabla_x(u(z) - u(\bar{z})) \rangle|$$

(since $u \in C_B^{1,\alpha}(\Omega_0)$ and by (2.22))

$$\leq c d_{\bar{z}}(\bar{z}, z)^\alpha d_{\bar{z}}(\bar{z}, \zeta).$$

□

3. PARAMETRICES

The proof of Theorem 1.3 is based on a representation formula for classical solutions to (1.1) in terms of the fundamental solution $\Gamma_{\bar{z}}$ of the frozen operator $L_{\bar{z}}$ in (2.4). In this section, we provide some crucial estimates of $\Gamma_{\bar{z}}$, with $\bar{z} \in \Omega_0 = \Omega \times \mathbb{R}^{N-1}$ (cf. Proposition 3.1). Most of the results of this section are rather technical.

We denote by $\Gamma_{\bar{z}}(z, \zeta)$ (resp. $\Gamma_H(\theta)$) the fundamental solution of $L_{\bar{z}}$ (resp. of L_H in (2.9)), evaluated in z (resp. in θ) and with pole in ζ (resp. in the origin). We note that

$$(3.1) \quad \Gamma_{\bar{z}}(z, \zeta) = \Gamma_{\bar{z}}(\zeta^{-1} \circ z, 0) = \frac{1}{\partial_{x_1} b(\bar{z})} \Gamma_H(\theta_{\bar{z}}^{\bar{z}}(\zeta^{-1} \circ z)),$$

where the product “ $\circ$ ” is defined in (2.11).

We introduce some auxiliary notation. We denote the identity by $D_0 = D_0^{\bar{z}} = D_0^H$ and for every multi-index $\sigma = (\sigma_1, \dots, \sigma_m) \in \{0, 1, \dots, 2N+1\}^m$, we set

$$D_{\sigma}^{\bar{z}} = D_{\sigma_1}^{\bar{z}} \cdots D_{\sigma_m}^{\bar{z}}, \quad D_{\sigma}^H = D_{\sigma_1}^H \cdots D_{\sigma_m}^H.$$

We call the weight of σ the number

$$(3.2) \quad |\sigma| = m_1^{\sigma} + 2m_2^{\sigma} + 3m_3^{\sigma}$$

where $m_1^\sigma, m_2^\sigma, m_3^\sigma$, respectively, are the cardinalities of the sets

$$\begin{aligned} &\{\sigma_j \in \sigma \mid 1 \leq \sigma_j \leq N\}, \\ &\{\sigma_j \in \sigma \mid \sigma_j = N+1\}, \\ &\{\sigma_j \in \sigma \mid N+2 \leq \sigma_j \leq 2N+1\}. \end{aligned}$$

Analogously, if $b \in C_b^{|\sigma|-2, \alpha}(\Omega)$, we put

$$D_\sigma = D_{\sigma_1} \cdots D_{\sigma_m}.$$

For greater convenience, whenever we consider a derivative D_σ with $|\sigma| \geq 2$, in the sequel we agree to assume $b \in C_b^{|\sigma|-2, \alpha}(\Omega)$. Moreover, to avoid any ambiguity, when we have a function F which depends on several variables, we systematically write $D(z)F$ instead of $DF(z)$. Then, for instance, $D(z)\Gamma_{\bar{z}}(\cdot, \zeta)$ denotes the D -derivative of $\Gamma_{\bar{z}}(\cdot, \zeta)$ evaluated at the point z .

The following estimate of $\Gamma_{\bar{z}}$ and its derivatives is well known: for every $\bar{z} \in \Omega_0$ and multi-index $\sigma \in \{0, 1, \dots, 2N+1\}^m$ there exists a positive constant c , such that

$$(3.3) \quad |D_\sigma^\bar{z}(z)\Gamma_{\bar{z}}(\cdot, \zeta)| \leq c d_{\bar{z}}(z, \zeta)^{-Q+2-|\sigma|}, \quad \forall z, \zeta \in \Omega_0, \quad z \neq \zeta.$$

Moreover, the constant c in (3.3) depends continuously on $\bar{z}$.

In the proof of Theorem 1.3, we shall also need to compare the fundamental solutions of frozen operators corresponding to different points of Ω_0 . The main result of this section is the following.

Proposition 3.1. *Let $\sigma \in \{1, \dots, N+1\}^m$. There exists a positive constant $c = c(\sigma, \Omega_0)$ such that*

$$(3.4) \quad |D_\sigma(z)\Gamma_{\bar{z}}(\cdot, \zeta) - D_\sigma(\bar{z})\Gamma_{\bar{z}}(\cdot, \zeta)| \leq c d_{\bar{z}}(\bar{z}, \zeta)^{-Q+2-|\sigma|-\alpha} d_{\bar{z}}(\bar{z}, z)^\alpha,$$

and, if $|\sigma| \geq 3$,

$$(3.5) \quad |D_\sigma(z)\Gamma_z(\cdot, \zeta) - D_\sigma(\bar{z})\Gamma_{\bar{z}}(\cdot, \zeta)| \leq c d_{\bar{z}}(\bar{z}, \zeta)^{-Q+2-|\sigma|-\alpha} d_{\bar{z}}(\bar{z}, z)^\alpha,$$

for every $z, \bar{z} \in \Omega_0$ and $\zeta \in \Omega_M(\bar{z}, z)$ (cf. (2.21)).

The proof of Proposition 3.1 is based on two lemmas. The first one gives an expression of D_j in terms of the frozen vector fields in (2.2).

Lemma 3.2. *Let $z, \bar{z} \in \Omega_0$ and $\sigma \in \{1, \dots, N+1\}^m$ (and $b \in C_b^{|\sigma|-2, \alpha}(\Omega)$, if $|\sigma| \geq 2$). For every smooth function φ , we have*

$$(3.6) \quad D_\sigma \varphi(z) = \sum_{\mu \in \{1, \dots, N+2\}^m} \Lambda_{\mu, \bar{z}}(z) (R_{\bar{z}}(z))^{\alpha_\mu} D_\mu^\bar{z} \varphi(z)$$

where $|\mu| \leq |\sigma| + \alpha_\mu$, $\alpha_\mu \leq m_2^\sigma$, and

$$(3.7) \quad \begin{aligned} R_{\bar{z}}(z) &= \frac{b(z) - P_{\bar{z}}^1 b(z)}{\partial_{x_1} b(\bar{z})} = \frac{b(z) - b(\bar{z}) - \langle \nabla_x b(\bar{z}), x - \bar{x} \rangle}{\partial_{x_1} b(\bar{z})}, \\ \Lambda_{\mu, \bar{z}}(z) &= \sum_i c_i(\mu) \prod_{\nu \in J_{\mu, i}} (D_\nu(z) R_{\bar{z}})^{\beta_\nu}. \end{aligned}$$

In (3.7), $J_{\mu, i}$ is a suitable subset of $\{1, \dots, N+1\}^{m-1}$, $|\nu| \leq |\sigma| - 2$ and $\beta_\nu \leq m$. Moreover, we have

$$(3.8) \quad |\Lambda_{\mu, \bar{z}}(z) (R_{\bar{z}}(z))^{\alpha_\mu} - \Lambda_{\mu, \bar{z}}(\bar{z}) (R_{\bar{z}}(\bar{z}))^{\alpha_\mu}| \leq c d_{\bar{z}}(\bar{z}, z)^{\alpha_\mu + \alpha}$$

for some positive constant $c = c(\sigma, \Omega_0)$.

Proof. We proceed by induction on $|\sigma|$. If $|\sigma| = 1, 2$, the assertion is trivial since $D_\sigma = D_\sigma^z$ if $m_\sigma^z = 0$ (cf. (3.2)), and we have

$$D_{N+1} = D_{N+1}^z + R_{\bar{z}} D_{N+2}^z.$$

We next consider $|\sigma| \geq 2$. If $j = 1, \dots, N$, by induction, we have

$$\begin{aligned} D_j D_\sigma \varphi(z) &= \sum_{\mu \in \{1, \dots, N+2\}^m} (D_j(z) (\Lambda_{\mu, \bar{z}} (R_{\bar{z}})^{\alpha_\mu}) D_\mu^z \varphi(z) \\ &\quad + \Lambda_{\mu, \bar{z}}(z) (R_{\bar{z}}(z))^{\alpha_\mu} D_j^z D_\mu^z \varphi(z)). \end{aligned}$$

Analogously, we have

$$\begin{aligned} D_{N+1} D_\sigma \varphi(z) &= \sum_{\mu \in \{1, \dots, N+2\}^m} \left(D_{N+1}(z) (\Lambda_{\mu, \bar{z}} (R_{\bar{z}})^{\alpha_\mu}) D_\mu^z \varphi(z) \right. \\ &\quad \left. + \Lambda_{\mu, \bar{z}}(z) (R_{\bar{z}}(z))^{\alpha_\mu} D_{N+1}^z D_\mu^z \varphi(z) + \Lambda_{\mu, \bar{z}}(z) (R_{\bar{z}}(z))^{\alpha_\mu+1} D_{N+2}^z D_\mu^z \varphi(z) \right). \end{aligned}$$

Then the thesis is a straightforward verification. In particular, (3.8) is a consequence of the fact that, by assumption, $b \in C_b^{|\sigma|-2, \alpha}(\Omega)$ and $R_{\bar{z}}(\bar{z}) = 0$. $\square$

The proof of (3.5) in Proposition 3.1 is rather delicate since we need to estimate the fundamental solutions of frozen operators related to different points $z, \bar{z}$. Here we use the canonical change of coordinates (2.5) and we investigate the properties of the fundamental solution Γ_H . The next lemma provides us with some basic estimates.

Lemma 3.3. *Let $b \in C_b^{1, \alpha}(\Omega)$. There exists a positive constant $c = c(\Omega)$, such that*

$$(3.9) \quad d_H(\theta_{\bar{z}}^z(\zeta), \theta_z^z(\zeta)) \leq c d_{\bar{z}}(\bar{z}, \zeta)^{1-\frac{\alpha}{3}} d_{\bar{z}}(\bar{z}, z)^{\frac{\alpha}{3}},$$

for every $z, \bar{z} \in \Omega_0$ and $\zeta \in \Omega_M(\bar{z}, z)$.

Proof. Let us denote $(-\theta_z^z(\zeta)) \oplus \theta_{\bar{z}}^z(\zeta) = \theta$. Keeping in mind formulas (2.5) and (2.6), we get

$$(3.10) \quad d_H(\theta_{\bar{z}}^z(\zeta), \theta_z^z(\zeta)) = |x - \bar{x}| + |t - \bar{t}|^{\frac{1}{2}} + \sum_{j=N+2}^{2N+1} |\theta_j|^{\frac{1}{3}}.$$

If $2 \leq j \leq N$, we have

$$\begin{aligned} \theta_{N+1+j} &= \eta_j - \bar{\eta}_j + \frac{(\tau - \bar{t})(x_j + \bar{x}_j)}{2} - \left(\eta_j - y_j + \frac{(\tau - t)(\xi_j + x_j)}{2} \right) \\ &\quad + \frac{(\xi_j - x_j)(\tau - \bar{t})}{2} - \frac{(\xi_j - x_j)(\tau - t)}{2} \\ &= \left(y_j - \bar{y}_j + \frac{(t - \bar{t})(x_j + \bar{x}_j)}{2} \right) + (\xi_j - x_j)(\tau - \bar{t}) - (\xi_j - \bar{x}_j)(\tau - t); \end{aligned}$$

therefore,

$$\begin{aligned} |\theta_{N+1+j}|^{\frac{1}{3}} &\leq d_{\bar{z}}(\bar{z}, z) + d_{\bar{z}}(\bar{z}, \zeta)^{\frac{2}{3}} d_{\bar{z}}(\bar{z}, z)^{\frac{1}{3}} + d_z(z, \zeta)^{\frac{2}{3}} d_{\bar{z}}(\bar{z}, z)^{\frac{1}{3}} \\ (3.11) \quad &\quad \text{(since } \zeta \in \Omega_M(\bar{z}, z) \text{ and by (2.22))} \\ &\leq c d_{\bar{z}}(\bar{z}, \zeta)^{1-\frac{\alpha}{3}} d_{\bar{z}}(\bar{z}, z)^{\frac{\alpha}{3}}. \end{aligned}$$

Next we prove the inequality

$$(3.12) \quad |\theta_{N+2}|^{\frac{1}{3}} \leq c d_{\bar{z}}(\bar{z}, \zeta)^{1-\frac{\alpha}{3}} d_{\bar{z}}(\bar{z}, z)^{\frac{\alpha}{3}},$$

which by (3.10) suffices to conclude the proof of the lemma. Noting that

$$\eta_j - y_j + x_j(\tau - t) = \eta_j - y_j + \frac{(\tau - t)(\xi_j + x_j)}{2} - \frac{(\tau - t)(\xi_j - x_j)}{2},$$

we have

$$(3.13) \quad \begin{aligned} \theta_{N+2} = & \frac{1}{2}(t - \bar{t})(\xi_1 - x_1 + \xi_1 - \bar{x}_1) + \frac{\eta - \bar{y} + (\tau - \bar{t})b(\bar{z})}{\partial_{x_1} b(\bar{z})} - \frac{\eta - y + (\tau - t)b(z)}{\partial_{x_1} b(z)} \\ & - \sum_{j=2}^N \left[\frac{\partial_{x_j} b(\bar{z})}{\partial_{x_1} b(\bar{z})} (\eta_j - \bar{y}_j + \bar{x}_j(\tau - \bar{t})) - \frac{\partial_{x_j} b(z)}{\partial_{x_1} b(z)} (\eta_j - y_j + x_j(\tau - t)) \right] \\ = & \frac{1}{2}(t - \bar{t})(\xi_1 - x_1 + \xi_1 - \bar{x}_1) + A_1 + A_2 + A_3 + A_4, \end{aligned}$$

where

$$\begin{aligned} A_1 = & \frac{y - \bar{y} + (t - \bar{t})b(\bar{z})}{\partial_{x_1} b(\bar{z})} + \frac{(\tau - t)(b(z) - b(\bar{z}))}{\partial_{x_1} b(\bar{z})} \\ & + \frac{\partial_{x_1} b(z) - \partial_{x_1} b(\bar{z})}{\partial_{x_1} b(\bar{z}) \partial_{x_1} b(z)} (\eta - y + (\tau - t)b(z)), \end{aligned}$$

so that, since $b \in C_b^{1,\alpha}(\Omega) \cap C^1(\Omega)$, we have

$$(3.14) \quad \begin{aligned} |A_1|^{\frac{1}{3}} & \leq c \left(d_{\bar{z}}(\bar{z}, z) + d_z(z, \zeta)^{\frac{2}{3}} d_{\bar{z}}(\bar{z}, z)^{\frac{1}{3}} + d_{\bar{z}}(\bar{z}, z)^{\frac{\alpha}{3}} d_z(z, \zeta) \right) \\ & \quad (\text{since } \zeta \in \Omega_M(\bar{z}, z) \text{ and by (2.22)}) \\ & \leq c d_{\bar{z}}(\bar{z}, \zeta)^{1-\frac{\alpha}{3}} d_{\bar{z}}(\bar{z}, z)^{\frac{\alpha}{3}}. \end{aligned}$$

Also,

$$A_2 = \sum_{j=2}^N \left(\frac{\partial_{x_j} b(z)}{\partial_{x_1} b(z)} - \frac{\partial_{x_j} b(\bar{z})}{\partial_{x_1} b(\bar{z})} \right) (\eta_j - \bar{y}_j + \bar{x}_j(\tau - \bar{t}))$$

so that, since $b \in C_b^{1,\alpha}(\Omega)$, we have

$$(3.15) \quad |A_2|^{\frac{1}{3}} \leq c d_{\bar{z}}(\bar{z}, z)^{\frac{\alpha}{3}} d_{\bar{z}}(\bar{z}, \zeta).$$

Moreover,

$$A_3 = \sum_{j=2}^N \frac{\partial_{x_j} b(z)}{\partial_{x_1} b(z)} \left(\eta_j - y_j + \frac{(\tau - t)(x_j + x_j)}{2} - \left(\eta_j - \bar{y}_j + \frac{(\tau - \bar{t})(\xi_j + \bar{x}_j)}{2} \right) \right)$$

which can be estimated as before:

$$(3.16) \quad |A_3|^{\frac{1}{3}} \leq c d_{\bar{z}}(\bar{z}, \zeta)^{1-\frac{\alpha}{3}} d_{\bar{z}}(\bar{z}, z)^{\frac{\alpha}{3}}.$$

Finally,

$$\begin{aligned} A_4 = & \frac{1}{2} \sum_{j=2}^N \frac{\partial_{x_j} b(z)}{\partial_{x_1} b(z)} ((\tau - \bar{t})(\xi_j - \bar{x}_j) - (\tau - t)(\xi_j - x_j)) \\ = & \frac{1}{2} \sum_{j=2}^N \frac{\partial_{x_j} b(z)}{\partial_{x_1} b(z)} ((t - \bar{t})(\xi_j - \bar{x}_j) + (\tau - t)(x_j - \bar{x}_j)) \end{aligned}$$

so that

$$\begin{aligned}
 |A_4|^{\frac{1}{3}} &\leq c d_{\bar{z}}(\bar{z}, z)^{\frac{2}{3}} d_{\bar{z}}(\bar{z}, \zeta)^{\frac{1}{3}} + d_z(z, \zeta)^{\frac{2}{3}} d_{\bar{z}}(\bar{z}, z)^{\frac{1}{3}} \\
 (3.17) \quad &\quad \quad \quad (\text{since } \zeta \in \Omega_M(\bar{z}, z) \text{ and by (2.22)}) \\
 &\leq c d_{\bar{z}}(\bar{z}, \zeta)^{1-\frac{\alpha}{3}} d_{\bar{z}}(\bar{z}, z)^{\frac{\alpha}{3}}.
 \end{aligned}$$

Plugging inequalities (3.14), (3.15), (3.16) and (3.17) back into (3.13), we obtain (3.12). This concludes the proof. $\square$

Remark 3.4. Let $b \in C_b^{1,\alpha}(\Omega)$. If $z, \bar{z} \in \Omega_0$, $\zeta \in \Omega_M(\bar{z}, z)$ and $\theta \in \mathbb{R}^{N+1}$ are such that

$$d_H(\theta_{\bar{z}}^{\bar{z}}(\zeta), \theta) \leq d_H(\theta_{\bar{z}}^{\bar{z}}(\zeta), \theta_z^z(\zeta)),$$

then there exists a positive constant $c = c(\Omega_0)$ such that

$$(3.18) \quad \|\theta\|_H \geq c \|\theta_{\bar{z}}^{\bar{z}}(\zeta)\|_H.$$

Indeed, we have

$$\begin{aligned}
 d_{\bar{z}}(\bar{z}, \zeta) &= \|\theta_{\bar{z}}^{\bar{z}}(\zeta)\|_H \leq c (d_H(0, \theta) + d_H(\theta_{\bar{z}}^{\bar{z}}(\zeta), \theta)) \\
 &\leq (d_H(0, \theta) + d_H(\theta_{\bar{z}}^{\bar{z}}(\zeta), \theta_z^z(\zeta)))
 \end{aligned}$$

(by (3.9) and since $\zeta \in \Omega_M(\bar{z}, z)$)

$$\leq c (d_H(0, \theta) + M^{-\frac{\alpha}{3}} d_{\bar{z}}(\bar{z}, \zeta)).$$

Thus, if M is suitably large, we get (3.18).

We are now in position to prove Proposition 3.1.

Proof of Proposition 3.1. We first prove estimate (3.4) by using Lemma 3.2. We have

$$|(D_\sigma(z) - D_\sigma(\bar{z})) \Gamma_{\bar{z}}(\cdot, \zeta)| \leq A_1 + A_2,$$

where

$$A_1 = \sum_{\mu \in \{1, \dots, N+2\}^m} |\Lambda_{\mu, \bar{z}}(z) (R_{\bar{z}}(z))^{\alpha_\mu} - \Lambda_{\mu, \bar{z}}(\bar{z}) (R_{\bar{z}}(\bar{z}))^{\alpha_\mu}| |D_{\mu}^{\bar{z}}(\bar{z}) \Gamma_{\bar{z}}(\cdot, \zeta)|$$

and

$$A_2 = \sum_{\mu \in \{1, \dots, N+2\}^m} |\Lambda_{\mu, \bar{z}}(z) (R_{\bar{z}}(z))^{\alpha_\mu} (D_{\mu}^{\bar{z}}(z) - D_{\mu}^{\bar{z}}(\bar{z})) \Gamma_{\bar{z}}(\cdot, \zeta)|.$$

Hence, by (3.3) and (3.8), we have

$$A_1 \leq c \sum_{\mu \in \{1, \dots, N+2\}^m} d_{\bar{z}}(\bar{z}, z)^{\alpha_\mu + \alpha} d_{\bar{z}}(\bar{z}, \zeta)^{-Q+2-|\mu|}$$

(since $\zeta \in \Omega_M(\bar{z}, z)$ and, by Lemma (3.2), $|\mu| \leq |\sigma| + \alpha_\mu$)

$$\leq c d_{\bar{z}}(\bar{z}, z)^\alpha d_{\bar{z}}(\bar{z}, \zeta)^{-Q+2-|\sigma|}.$$

On the other hand, by the mean value theorem, for some $\tilde{z}$ such that $d_{\tilde{z}}(\tilde{z}, \tilde{z}) \leq d_{\tilde{z}}(\tilde{z}, z)$, we have

$$\begin{aligned} |(D_{\mu}^{\tilde{z}}(z) - D_{\mu}^{\tilde{z}}(\tilde{z})) \Gamma_{\tilde{z}}(\cdot, \zeta)| &= |\langle \theta_{\tilde{z}}^{\tilde{z}}(z), \nabla_{\tilde{z}} D_{\mu}^{\tilde{z}}(\tilde{z}) \Gamma_{\tilde{z}}(\cdot, \zeta) \rangle| \\ &\leq c \sum_{j=1}^3 d_{\tilde{z}}(\tilde{z}, z)^j d_{\tilde{z}}(\tilde{z}, \zeta)^{-Q+2-|\mu|-j} \end{aligned}$$

(since $\zeta \in \Omega_M(\tilde{z}, z)$)

$$\leq c d_{\tilde{z}}(\tilde{z}, \zeta)^{-Q+2-|\mu|-\alpha} d_{\tilde{z}}(\tilde{z}, z)^{\alpha}.$$

As before, we conclude that

$$A_2 \leq c d_{\tilde{z}}(\tilde{z}, \zeta)^{-Q+2-|\sigma|-\alpha} d_{\tilde{z}}(\tilde{z}, z)^{\alpha}.$$

Next we prove (3.5). By (3.1) and Lemma 3.2, we have

$$\begin{aligned} &|D_{\sigma}(z) \Gamma_z(\cdot, \zeta) - D_{\sigma}(\tilde{z}) \Gamma_{\tilde{z}}(\cdot, \zeta)| \\ &\leq \left| \sum_{|\mu| \leq |\sigma|} \left(\frac{\Lambda_{\mu, z}(z)}{\partial_{x_1} b(z)} D_{\mu}^H(-\theta_z^z(\zeta)) - \frac{\Lambda_{\mu, \tilde{z}}(\tilde{z})}{\partial_{x_1} b(\tilde{z})} D_{\mu}^H(-\theta_{\tilde{z}}^{\tilde{z}}(\zeta)) \right) \Gamma_H \right| \leq A_1 + A_2 \end{aligned}$$

where

$$\begin{aligned} A_1 &= \sum_{|\mu| \leq |\sigma|} \left| \frac{\Lambda_{\mu, z}(z)}{\partial_{x_1} b(z)} - \frac{\Lambda_{\mu, \tilde{z}}(\tilde{z})}{\partial_{x_1} b(\tilde{z})} \right| |D_{\mu}^H \Gamma_H(-\theta_z^z(\zeta))|, \\ A_2 &= \sum_{|\mu| \leq |\sigma|} \left| \frac{\Lambda_{\mu, \tilde{z}}(\tilde{z})}{\partial_{x_1} b(\tilde{z})} \right| |D_{\mu}^H \Gamma_H(-\theta_z^z(\zeta)) - D_{\mu}^H \Gamma_H(-\theta_{\tilde{z}}^{\tilde{z}}(\zeta))|. \end{aligned}$$

Since the function

$$z \mapsto \frac{\Lambda_{\mu, z}(z)}{\partial_{x_1} b(z)} \in C_B^{\alpha}(\Omega_0),$$

we get

$$A_1 \leq c d_{\tilde{z}}(\tilde{z}, z)^{\alpha} d_z(z, \zeta)^{-Q+2-|\sigma|}.$$

On the other hand, by the mean value theorem, there exists θ such that

$$d_H(\theta_{\tilde{z}}^{\tilde{z}}(\zeta), \theta) \leq d_H(\theta_{\tilde{z}}^{\tilde{z}}(\zeta), \theta_z^z(\zeta)),$$

and

$$\begin{aligned} |D_{\mu}^H \Gamma_H(-\theta_z^z(\zeta)) - D_{\mu}^H \Gamma_H(-\theta_{\tilde{z}}^{\tilde{z}}(\zeta))| &= |\langle (-\theta_z^z(\zeta)) \oplus \theta_{\tilde{z}}^{\tilde{z}}(\zeta), \nabla_H D_{\mu}^H \Gamma_H(\theta) \rangle| \\ &\leq c \sum_{j=1}^3 d_H(\theta_{\tilde{z}}^{\tilde{z}}(\zeta), \theta_z^z(\zeta))^j \|\theta\|_H^{-Q+2-|\mu|-j} \\ &\leq c d_{\tilde{z}}(\tilde{z}, z)^{\alpha} d_z(z, \zeta)^{-Q+2-|\sigma|-\alpha}, \end{aligned}$$

where the last inequality follows from Lemma 3.3 and Remark 3.4 (note that the assumption $b \in C_b^{1, \alpha}(\Omega)$ of Lemma 3.3 is fulfilled because $|\sigma| \geq 3$). The proof is completed. $\square$

We end this section by describing the fine behaviour of the vector fields $D_j^{\tilde{z}}$ through the right translations w.r.t. the law “o” in (2.11).

Lemma 3.5. *Let $b \in C_b^{1,\alpha}(\Omega)$. For every smooth function φ and multi-index $\mu \in \{1, \dots, N+2\}^m$, we have*

$$(3.19) \quad D_\mu^{\bar{z}}(z)\varphi(\cdot \circ \zeta^{-1}) = \sum_{\nu \in \{1, \dots, 2N+1\}^m} P_{\nu, \bar{z}}(\zeta) D_\nu^{\bar{z}}(z \circ \zeta^{-1})\varphi,$$

where $P_{\nu, \bar{z}}$ is a polynomial $\delta_\lambda^{\bar{z}}$ -homogeneous of degree k_ν , with $|\nu| \leq |\mu| + k_\nu$, whose coefficients are functions of $\bar{z}$ of class $C_B^\alpha(\Omega_0)$.

Proof. Since

$$z \circ \zeta^{-1} = (x - \xi, y - \eta + \mathcal{J}_x B(\bar{z})(\tau(x - \bar{x}) - (\tau - \bar{t})\xi), t - \tau),$$

we have, for $2 \leq j \leq N$,

$$D_1^{\bar{z}}(z)\varphi(\cdot \circ \zeta^{-1}) = D_1^{\bar{z}}(z \circ \zeta^{-1})\varphi + \tau D_{N+2}^{\bar{z}}(z \circ \zeta^{-1})\varphi,$$

$$D_j^{\bar{z}}(z)\varphi(\cdot \circ \zeta^{-1}) = D_j^{\bar{z}}(z \circ \zeta^{-1})\varphi + \tau \left(D_{N+1+j}^{\bar{z}} - \frac{1}{\partial_{x_1} b(\bar{z})} D_{N+2}^{\bar{z}} \right) (z \circ \zeta^{-1})\varphi,$$

$$D_{N+1}^{\bar{z}}(z)\varphi(\cdot \circ \zeta^{-1}) = \left(D_{N+1}^{\bar{z}} + \sum_{j=2}^N \xi_j D_{N+1+j}^{\bar{z}} \right) (z \circ \zeta^{-1})\varphi,$$

$$D_{N+2}^{\bar{z}}(z)\varphi(\cdot \circ \zeta^{-1}) = D_{N+2}^{\bar{z}}(z \circ \zeta^{-1})\varphi.$$

The thesis easily follows by induction on $|\mu|$. $\square$

4. HÖLDER REGULARITY

In this section we prove Theorem 1.3. We first give the definition of a classical solution of (1.1).

Definition 4.1. A classical solution of (1.1) is a function $u \in C^1(\Omega)$, with second-order derivatives $\partial_{x_i x_j} u \in C(\Omega)$, $1 \leq i, j \leq N$, verifying (1.1).

For greater convenience, since in the following proof we deal with several estimates, we shall denote by c a constant which will not be always the same.

Proof of Theorem 1.3. We proceed by induction on k . Let us remark that, if $k \geq 7$, the proof is standard and the thesis follows by differentiating the equation. More precisely, if $b, f \in C_b^{k-2, \alpha}(\Omega)$, with $0 < \alpha < 1$ and $k \geq 7$, then by the inductive hypothesis, $u \in C_b^{k-1, \bar{\alpha}}(\Omega)$ for every $\bar{\alpha} \in]0, \alpha[$. Subsequently, by differentiating equation (1.1) w.r.t. the variable y and by denoting, as usual, $\partial_y u = u_y$, we get

$$L_b u_y = f_y - b_y u_y \in C_b^{k-5, \bar{\alpha}}(\Omega), \quad \forall \bar{\alpha} \in]0, \alpha[.$$

Thus, we deduce that $u_y \in C_b^{k-3, \bar{\alpha}}(\Omega)$. Next we differentiate equation (1.1) w.r.t. the variable x_j , $j = 1, \dots, N-1$, and we get

$$L_b u_{x_j} = f_{x_j} - b_{x_j} u_y \in C_b^{k-3, \bar{\alpha}}(\Omega), \quad \forall \bar{\alpha} \in]0, \alpha[.$$

Therefore, we deduce that $u_{x_j} \in C_b^{k-1, \bar{\alpha}}(\Omega)$. Finally, we differentiate equation (1.1) w.r.t. to $X_{N+1} = b\partial_y - \partial_t$:

$$L_b(X_{N+1}u) = X_{N+1}f + 2\langle \nabla_x b, \nabla_x u_y \rangle + u_y \Delta_x b \in C_b^{k-4, \bar{\alpha}}(\Omega), \quad \forall \bar{\alpha} \in]0, \alpha[.$$

Hence, $X_{N+1}u \in C_b^{k-2, \bar{\alpha}}(\Omega)$ and this proves that $u \in C_b^{k, \bar{\alpha}}(\Omega)$ for every $\bar{\alpha} \in]0, \alpha[$.

We next consider $3 \leq k \leq 6$. We set the problem in dimension $2N+1$ and we prove that $u \in C_B^{k, \bar{\alpha}}(\Omega_0)$. Then, by Remark 2.5, we infer that $u \in C_b^{k, \bar{\alpha}}(\Omega)$. We split

the proof into two steps: existence of the derivatives and Hölder continuity. Since the case $k = 2$ is easier, we shall only sketch its proof separately at the end. $\square$

The case $3 \leq k \leq 6$: existence of the derivatives. Let $b, f \in C_b^{k-2, \alpha}(\Omega)$, with $0 < \alpha < 1$ and $3 \leq k \leq 6$. By the inductive hypothesis, $u \in C_b^{k-1, \bar{\alpha}}(\Omega)$ for every $\bar{\alpha} \in]0, \alpha[$. Since we aim to prove a local result, we only prove the existence of the derivatives of order k of u in a domain E_0 , where $\bar{E}_0$ is a compact subset of Ω_0 . To this end, we represent the solution u in terms of the fundamental solution $\Gamma_{\bar{z}}$, $\bar{z} \in E_0$. More precisely, we consider a cut-off function $\varphi \in C_0^\infty(\Omega_0)$ such that $\varphi = 1$ in a neighborhood of E_0 . We remark that it is not restrictive to assume that

$$(4.1) \quad \text{supp}(\nabla \varphi) \subseteq \Omega_M(\bar{z}, z), \quad \forall z, \bar{z} \in E_0,$$

where M is large so that (2.22) holds for every $z, \bar{z} \in E_0$. We have

$$(4.2) \quad \begin{aligned} (u\varphi)(z) &= \int_{\Omega_0} \Gamma_{\bar{z}}(z, \zeta) L_{\bar{z}}(u\varphi)(\zeta) d\zeta \\ &= \int_{\Omega_0} \Gamma_{\bar{z}}(z, \zeta) (\varphi (f - (b - P_{\bar{z}}^1 b) u_{y_1}) + 2\langle \nabla_x u, \nabla_x \varphi \rangle + u L_{\bar{z}} \varphi) d\zeta. \end{aligned}$$

Consequently, the solution u can be expressed in the form

$$u(z) = I_{1, \bar{z}}(z) - I_{2, \bar{z}}(z), \quad z \in E_0$$

where

$$I_{j, \bar{z}}(z) = \int_{\Omega_0} \Gamma_{\bar{z}}(z, \zeta) U_{j, \bar{z}}(\zeta) d\zeta, \quad j = 1, 2,$$

with $U_{1, \bar{z}} \in C_0^\infty(\Omega_0)$ and

$$(4.3) \quad |U_{2, \bar{z}}(\zeta)| \leq c d_{\bar{z}}(\bar{z}, \zeta)^{k-2+\bar{\alpha}}, \quad \forall \bar{z} \in E_0, \zeta \in \Omega_0, \bar{\alpha} \in]0, \alpha[.$$

Indeed, it suffices to put

$$(4.4) \quad \begin{aligned} U_{1, \bar{z}} &= \varphi (P_{\bar{z}}^{k-2} f + P_{\bar{z}}^{k-4} u_{y_1} (P_{\bar{z}}^{k-2} b - P_{\bar{z}}^1 b)) \\ &\quad + 2\langle P_{\bar{z}}^{k-2} \nabla_x u, \nabla_x \varphi \rangle + (L_{\bar{z}} \varphi) P_{\bar{z}}^{k-2} u, \end{aligned}$$

$$(4.5) \quad \begin{aligned} U_{2, \bar{z}} &= \varphi (u_{y_1} (b - P_{\bar{z}}^{k-2} b) + (P_{\bar{z}}^{k-2} b - P_{\bar{z}}^1 b) (u_{y_1} - P_{\bar{z}}^{k-4} u_{y_1}) + f - P_{\bar{z}}^{k-2} f) \\ &\quad + 2\langle \nabla_x u - P_{\bar{z}}^{k-2} \nabla_x u, \nabla_x \varphi \rangle + L_{\bar{z}} \varphi (u - P_{\bar{z}}^{k-2} u) \end{aligned}$$

where we agree that $P_{\bar{z}}^h u \equiv 0$ if $h < 0$. Hence, (4.3) is a consequence of the regularity assumptions on u, b, f and of the estimate

$$P_{\bar{z}}^{k-2} b(\zeta) - P_{\bar{z}}^1 b(\zeta) = O(d_{\bar{z}}(\bar{z}, \zeta)^2), \quad \text{as } \zeta \longrightarrow \bar{z},$$

which can be easily deduced from the homogeneity property of the Taylor polynomial $P_{\bar{z}}^h b$. Note that $I_{1, \bar{z}}$ is a smooth function on E_0 . Indeed, by the change of variables $\bar{\zeta} = \zeta^{-1} \circ z$, we have

$$(4.6) \quad I_{1, \bar{z}}(z) = \int \Gamma_{\bar{z}}(\bar{\zeta}, 0) U_{1, \bar{z}}(z \circ \bar{\zeta}^{-1}) d\bar{\zeta}$$

and $U_{1, \bar{z}} \in C_0^\infty(\Omega_0)$.

On the other hand, the function

$$(4.7) \quad J_\sigma(\bar{z}) = \int_{\Omega_0} D_\sigma(\bar{z}) \Gamma_{\bar{z}}(\cdot, \zeta) U_{2, \bar{z}}(\zeta) d\zeta, \quad \bar{z} \in E_0,$$

is well defined for $|\sigma| = k$. Indeed, by Lemma 3.2, we have

$$(4.8) \quad |D_\sigma(\bar{z})\Gamma_{\bar{z}}(\cdot, \zeta)| = \left| \sum_{|\mu| \leq |\sigma|} \Lambda_{\mu, \bar{z}}(\bar{z}) D_\mu^{\bar{z}}(\bar{z}) \Gamma_{\bar{z}}(\cdot, \zeta) \right| \leq c d_{\bar{z}}(\bar{z}, \zeta)^{-Q+2-|\sigma|};$$

therefore, by (4.3), the integral in the right-hand side of (4.7) converges.

We claim that, for every multi-index $\sigma \in \{1, \dots, N+1\}^m$, with $|\sigma| \leq k$, we have

$$(4.9) \quad D_\sigma u(\bar{z}) = D_\sigma(\bar{z})I_{1, \bar{z}} - J_\sigma(\bar{z}), \quad \forall \bar{z} \in E_0.$$

The proof of (4.9) is based on the use of some suitable high-order difference quotients related to the system D_j , $j = 1, \dots, N+1$. For $z \in \Omega$ and $\delta \in \mathbb{R} \setminus \{0\}$ sufficiently small, we set

$$\begin{aligned} \Delta_{(j)}(u, z, \delta) &= \frac{u(\exp(\delta D_j)(z)) - u(z)}{\delta}, \quad 1 \leq j \leq N, \\ \Delta_{(N+1)}(u, z, \delta) &= \frac{u(\exp(\delta^2 D_{N+1})(z)) - u(z)}{\delta^2}. \end{aligned}$$

Also, if $\sigma \in \{1, \dots, N+1\}^m$, by recurrence we define

$$\Delta_\sigma(u, z, \delta) = \Delta_{(\sigma_m)}(\Delta_{(\sigma_1, \dots, \sigma_{m-1})}(u, \cdot, \delta), z, \delta).$$

The following result can be proved as in [6], Remark 4.2.

Lemma 4.2. *Let $u \in C_b^{k-1, \alpha}(\Omega)$ and $\sigma \in \{1, \dots, N+1\}^m$, $|\sigma| = k$. If there exists*

$$\lim_{\delta \rightarrow 0} \Delta_\sigma(u, \cdot, \delta) = v$$

uniformly on compact subsets of Ω , then there exists $D_\sigma u = v$.

We are now in a position to prove (4.9). Since $I_{1, \bar{z}}$ is a smooth function, by the mean value theorem, we have

$$\Delta_\sigma(I_{1, \bar{z}}, \bar{z}, \delta) = D_\sigma(z_\delta)I_{\bar{z}}^1$$

for some z_δ such that

$$(4.10) \quad d_{\bar{z}}(\bar{z}, z_\delta) \leq c_1 \delta,$$

where the constant c_1 depends only on σ and on the constant c_0 in (2.16) (for instance, $c_1 = m c_0^{m-1}$ is fine). Thus, $\Delta_\sigma(I_{1, \bar{z}}, \bar{z}, \delta)$ converges to $D_\sigma(\bar{z})I_{1, \bar{z}}$ as δ tends to zero uniformly in $\bar{z} \in E_0$. Therefore, by Lemma 4.2, in order to prove (4.9), it suffices to show that

$$(4.11) \quad \lim_{\delta \rightarrow 0} \Delta_\sigma(I_{2, \bar{z}}, \bar{z}, \delta) = J_\sigma(\bar{z})$$

uniformly on E_0 .

Let us consider a cut-off function $\chi \in C_0^\infty(\mathbb{R}, [0, 1])$ such that $\chi(s) = 1$ for $s \geq 2$ and χ vanishes for $s \leq 1$. We set

$$(4.12) \quad I_{2, \bar{z}, \delta}(z) = \int_{\Omega_0} \Gamma_{\bar{z}}(z, \zeta) \chi\left(\frac{d_{\bar{z}}(\bar{z}, \zeta)}{2c_1 M \delta}\right) U_{2, \bar{z}}(\zeta) d\zeta, \quad z \in E_0,$$

where M, c_1 are the constants in (4.1), (4.10) respectively. Note that

$$z \mapsto I_{2, \bar{z}, \delta}(z), \quad d_{\bar{z}}(\bar{z}, z) \leq c_1 \delta,$$

is a smooth function. Indeed, if $d_{\bar{z}}(\bar{z}, z) \leq c_1\delta$, then by (2.16), the argument of the cut-off function χ in (4.12) satisfies

$$\frac{d_{\bar{z}}(\bar{z}, \zeta)}{2c_1M\delta} \leq \frac{1}{2c_1\delta} (d_{\bar{z}}(\bar{z}, z) + d_z(z, \zeta)) \leq 1, \quad \forall \zeta, \quad d_z(z, \zeta) \leq c_1\delta,$$

so that χ vanishes in a neighborhood of the pole of $\Gamma_{\bar{z}}(z, \zeta)$.

We claim that

$$(4.13) \quad \sup_{d_{\bar{z}}(\bar{z}, z) \leq c_1\delta} |I_{2, \bar{z}, \delta}(z) - I_{2, \bar{z}}(z)| \leq c\delta^{k+\bar{\alpha}},$$

$$(4.14) \quad \sup_{d_{\bar{z}}(\bar{z}, z) \leq c_1\delta} |D_{\sigma}(z)I_{2, \bar{z}, \delta} - J_{\sigma}(\bar{z})| \leq c\delta^{\bar{\alpha}},$$

for some positive constant c which depends continuously on $\bar{z}$. Taking the claim for granted, we immediately conclude the proof of (4.11) (and consequently of (4.9)):

$$|\Delta_{\sigma}(I_{2, \bar{z}}, \bar{z}, \delta) - J_{\sigma}(\bar{z})|$$

(by the mean value theorem, for some z_{δ} such that $d_{\bar{z}}(\bar{z}, z_{\delta}) \leq c_1\delta$)

$$\leq |\Delta_{\sigma}(I_{2, \bar{z}}, \bar{z}, \delta) - \Delta_{\sigma}(I_{2, \bar{z}, \delta}, \bar{z}, \delta)| + |D_{\sigma}(z_{\delta})I_{2, \bar{z}, \delta} - J_{\sigma}(\bar{z})|$$

(by (4.13) and (4.14))

$$\leq c\delta^{\bar{\alpha}}.$$

We are left with the proof of the claim. We begin by proving (4.13). We have

$$|I_{2, \bar{z}}(z) - I_{2, \bar{z}, \delta}(z)| \leq \int_{d_{\bar{z}}(\bar{z}, \zeta) \leq 4c_1M\delta} \Gamma_{\bar{z}}(z, \zeta) |U_{2, \bar{z}}(\zeta)| d\zeta$$

(by (3.3) and (4.3))

$$\leq c\delta^{k-2+\bar{\alpha}} \int_{d_{\bar{z}}(\bar{z}, \zeta) \leq 4c_1M\delta} d_{\bar{z}}(z, \zeta)^{-Q+2} d\zeta$$

(since, by assumption, $d_{\bar{z}}(z, \zeta) \leq c_0(d_{\bar{z}}(\bar{z}, z) + d_{\bar{z}}(\bar{z}, \zeta)) \leq c_0(c_1\delta + d_{\bar{z}}(\bar{z}, \zeta))$)

$$\leq c\delta^{k-2+\bar{\alpha}} \int_{d_{\bar{z}}(z, \zeta) \leq c_0c_1(1+4M)\delta} d_{\bar{z}}(z, \zeta)^{-Q+2} d\zeta \leq c\delta^{k+\bar{\alpha}}.$$

Next, we prove (4.14). We have

$$|D_{\sigma}(z)I_{2, \bar{z}, \delta} - J_{\sigma}(\bar{z})| \leq A_1(\bar{z}, z) + A_2(\bar{z}, z),$$

where

$$A_1(\bar{z}, z) = \int_{\Omega_0} |D_{\sigma}(\bar{z})\Gamma_{\bar{z}}(\cdot, \zeta)| \left(1 - \chi\left(\frac{d_{\bar{z}}(\bar{z}, \zeta)}{2c_1M\delta}\right)\right) |U_{2, \bar{z}}(\zeta)| d\zeta$$

and

$$A_2(\bar{z}, z) = \int_{\Omega_0} |(D_{\sigma}(z) - D_{\sigma}(\bar{z}))\Gamma_{\bar{z}}(\cdot, \zeta)| \chi\left(\frac{d_{\bar{z}}(\bar{z}, \zeta)}{2c_1M\delta}\right) |U_{2, \bar{z}}(\zeta)| d\zeta.$$

Proceeding as in the proof of (4.13) and using (4.8), we obtain

$$A_1(\bar{z}, z) \leq c\delta^{\bar{\alpha}}.$$

On the other hand, we remark that, if $d_{\bar{z}}(\bar{z}, z) \leq c_1\delta$, then

$$\text{supp} \left(\chi \left(\frac{d_{\bar{z}}(\bar{z}, \cdot)}{2c_1M\delta} \right) \right) \subseteq \Omega_M.$$

Hence, by estimate (3.4) of Proposition 3.1 and (4.3), we get

$$A_2(\bar{z}, z) \leq c \int_{\Omega_M} d_{\bar{z}}(\bar{z}, z)^{\bar{\alpha}} d_{\bar{z}}(\bar{z}, \zeta)^{-Q+\bar{\alpha}} d\zeta \leq c\delta^{\bar{\alpha}}.$$

This concludes the proof of (4.14).

The case $3 \leq k \leq 6$: Hölder continuity of the derivatives. Let $\alpha \in]0, 1[$ and $3 \leq k \leq 6$. In the previous step, we have proved that, if $b, f \in C_b^{k-2, \alpha}(\Omega)$, and $u \in C_b^{k-1, \bar{\alpha}}(\Omega)$, for any $\bar{\alpha} \in]0, \alpha[$, then u is D_σ -differentiable for any σ of weight $|\sigma| \leq k$ and the representation formula (4.9) of $D_\sigma u$ holds.

In this step, we aim to prove that

$$(4.15) \quad D_\sigma u \in C_B^{\bar{\alpha}}(E_0), \quad \forall \sigma \in \{1, \dots, N+1\}^m, \quad |\sigma| = k, \text{ and } \bar{\alpha} \in]0, \alpha[,$$

$$(4.16) \quad D_\sigma u \in C_{D_{N+1}}^{1+\bar{\alpha}}(E_0), \quad \forall \sigma, \quad |\sigma| = k-1, \text{ and } \bar{\alpha} \in]0, \alpha[,$$

where E_0 is any domain fixed as in the preceding step. As a consequence of (4.15) and (4.16), we deduce that $u \in C_b^{k, \bar{\alpha}}(\Omega)$.

We prove (4.15) by means of Lemma 2.8 and we consider separately the terms $D_\sigma(\bar{z})I_{1, \bar{z}}$ and $J_\sigma(\bar{z})$ in formula (4.9), with $\bar{z} \in E_0$ and $|\sigma| = k$. Let us begin with J_σ and prove that

$$(4.17) \quad |J_\sigma(z) - J_\sigma(\bar{z})| \leq cd_{\bar{z}}(\bar{z}, z)^{\bar{\alpha}}, \quad \forall \bar{z}, z \in E_0, \quad \bar{\alpha} \in]0, \alpha[,$$

for some constant $c = c(E_0, \bar{\alpha}, k)$. By a standard decomposition, we obtain

$$J_\sigma(z) - J_\sigma(\bar{z}) = A_1(z, \bar{z}) + A_2(z, \bar{z}) + A_3(z, \bar{z}),$$

where

$$\begin{aligned} A_1(z, \bar{z}) &= \int_{\Omega_0 \setminus \Omega_M(\bar{z}, z)} (D_\sigma(z)\Gamma_z(\cdot, \zeta)U_{2, z}(\zeta) - D_\sigma(\bar{z})\Gamma_{\bar{z}}(\cdot, \zeta)U_{2, \bar{z}}(\zeta)) d\zeta, \\ A_2(z, \bar{z}) &= \int_{\Omega_M(\bar{z}, z)} (D_\sigma(z)\Gamma_z(\cdot, \zeta) - D_\sigma(\bar{z})\Gamma_{\bar{z}}(\cdot, \zeta)) U_{2, \bar{z}}(\zeta) d\zeta, \\ A_3(z, \bar{z}) &= \int_{\Omega_M(\bar{z}, z)} D_\sigma(z)\Gamma_z(\cdot, \zeta) (U_{2, z}(\zeta) - U_{2, \bar{z}}(\zeta)) d\zeta. \end{aligned}$$

By (4.3) and (4.8), it is straightforward to prove that

$$|A_1(z, \bar{z})| \leq cd_{\bar{z}}(\bar{z}, z)^{\bar{\alpha}}, \quad \forall \bar{z}, z \in E_0.$$

Next, we estimate $A_2(z, \bar{z})$ by using (3.5) of Proposition 3.1. We deduce that, for every $\bar{\alpha} \in]0, \alpha[$, there exists a constant $c = c(E_0, \bar{\alpha})$, such that

$$|A_2(z, \bar{z})| \leq c \int_{\Omega_M(\bar{z}, z)} d_{\bar{z}}(\bar{z}, \zeta)^{-Q} d_{\bar{z}}(\bar{z}, z)^{\bar{\alpha}} d\zeta \leq cd_{\bar{z}}(\bar{z}, z)^{\bar{\alpha}}, \quad \forall \bar{z}, z \in E_0.$$

Analogously, we have

$$(4.18) \quad |A_3(z, \bar{z})| \leq cd_{\bar{z}}(\bar{z}, z)^{\bar{\alpha}}, \quad \forall \bar{z}, z \in E_0.$$

Indeed, (4.18) is an immediate consequence of (4.8) and of the following estimate: there exists a positive constant $c = c(E_0)$ such that

$$(4.19) \quad |U_{2,z}(\zeta) - U_{2,\bar{z}}(\zeta)| \leq cd_{\bar{z}}(\bar{z}, z)^{\bar{\alpha}} d_{\bar{z}}(\bar{z}, \zeta)^{k-2}, \quad \forall z, \bar{z} \in E_0, \zeta \in \Omega_M(\bar{z}, z).$$

Let us prove (4.19). We have

$$U_{2,z} - U_{2,\bar{z}} = B_1(z, \bar{z}) + \varphi B_2(z, \bar{z}) + B_3(z, \bar{z}),$$

where

$$\begin{aligned} B_1(z, \bar{z}) &= \varphi (u_{y_1} (P_{\bar{z}}^{k-2} b - P_z^{k-2} b) + (P_{\bar{z}}^{k-2} f - P_z^{k-2} f)) \\ &\quad + 2(P_{\bar{z}}^{k-2} \nabla_x u - P_z^{k-2} \nabla_x u, \nabla_x \varphi), \\ B_2(z, \bar{z}) &= (P_z^{k-2} b - P_z^1 b) (u_{y_1} - P_z^{k-4} u_{y_1}) - (P_{\bar{z}}^{k-2} b - P_{\bar{z}}^1 b) (u_{y_1} - P_{\bar{z}}^{k-4} u_{y_1}), \\ B_3(z, \bar{z}) &= L_z \varphi (u - P_z^{k-2} u) - L_{\bar{z}} \varphi (u - P_{\bar{z}}^{k-2} u). \end{aligned}$$

By (2.23) of Lemma 2.10, we get at once

$$|B_1(z, \bar{z})(\zeta)| \leq cd_{\bar{z}}(\bar{z}, z)^{\bar{\alpha}} d_{\bar{z}}(\bar{z}, \zeta)^{k-2}, \quad \forall z, \bar{z} \in E_0, \zeta \in \Omega_M(\bar{z}, z).$$

We next remark that

$$\begin{aligned} B_2(z, \bar{z}) &= (P_z^{k-2} b - P_z^1 b) (P_{\bar{z}}^{k-4} u_{y_1} - P_z^{k-4} u_{y_1}) \\ &\quad + [(P_z^{k-2} b - P_z^1 b) - (P_{\bar{z}}^{k-2} b - P_{\bar{z}}^1 b)] (u_{y_1} - P_{\bar{z}}^{k-4} u_{y_1}). \end{aligned}$$

Therefore, by (2.23) and (2.24) of Lemma 2.10, we infer

$$|B_2(z, \bar{z})(\zeta)| \leq cd_{\bar{z}}(\bar{z}, z)^{\bar{\alpha}} d_{\bar{z}}(\bar{z}, \zeta)^{k-2}, \quad \forall z, \bar{z} \in E_0, \zeta \in \Omega_M(\bar{z}, z).$$

Finally, we observe that

$$\begin{aligned} B_3(z, \bar{z}) &= (L_z \varphi - L_{\bar{z}} \varphi) (u - P_z^{k-2} u) + L_{\bar{z}} \varphi (P_{\bar{z}}^{k-2} u - P_z^{k-2} u) \\ &= (P_z^1 b - P_{\bar{z}}^1 b) \partial_{y_1} \varphi (u - P_z^{k-2} u) + L_{\bar{z}} \varphi (P_{\bar{z}}^{k-2} u - P_z^{k-2} u). \end{aligned}$$

Thus, applying again Lemma 2.10, we establish (4.19) and, consequently, (4.17).

We conclude the proof of (4.15) by showing that

$$(4.20) \quad |D_\sigma(z)I_{1,z} - D_\sigma(\bar{z})I_{1,\bar{z}}| \leq cd_{\bar{z}}(\bar{z}, z)^\alpha, \quad \forall \bar{z}, z \in E_0,$$

for some constant $c = c(E_0, k)$. We denote $\bar{\zeta} = (\bar{\xi}, \bar{\eta}, \bar{\tau})$ and by

$$(4.21) \quad z \longmapsto R_{\zeta^{-1}}^{(\bar{z})}(z) \equiv z \circ \zeta^{-1} = (x - \xi, y - \eta + \mathcal{J}_x B(\bar{z})(\tau(x - \bar{x}) - (\tau - \bar{t})\xi), t - \tau),$$

the right translation w.r.t. the law “ $\circ$ ” in (2.11). By (4.6), for every $z \in E_0$ and multi-index $\sigma \in \{1, \dots, N+1\}^m$ of weight $|\sigma| = k$, we have

$$\begin{aligned}
 (4.22) \quad D_\sigma(z)I_{1,z} &= \int \Gamma_z(\bar{\zeta}, 0) D_\sigma(z) U_{1,z}(R_{\bar{\zeta}^{-1}}^{(z)}(\cdot)) d\bar{\zeta} \\
 &\quad \text{(by Lemma 3.2)} \\
 &= \sum_{|\mu| \leq k} \Lambda_{\mu,z}(z) \int \Gamma_z(\bar{\zeta}, 0) D_\mu^z(z) U_{1,z}(R_{\bar{\zeta}^{-1}}^{(z)}(\cdot)) d\bar{\zeta} \\
 &\quad \text{(by Lemma 3.5)} \\
 &= \sum_{|\mu| \leq k} \sum_{\nu \in \{1, \dots, 2N+1\}^m} \Lambda_{\mu,z}(z) \int \Gamma_z(\bar{\zeta}, 0) P_{\nu,z}(\bar{\zeta}) D_\nu^z(z \circ \bar{\zeta}^{-1}) U_{1,z} d\bar{\zeta} \\
 &\quad \text{(by the change of variables } \bar{\zeta} = \zeta^{-1} \circ z) \\
 &= \int_{\Omega_0} \Gamma_z(z, \zeta) V_z(\zeta) d\zeta,
 \end{aligned}$$

where

$$V_z(\zeta) = \sum_{|\mu| \leq k} \sum_{\nu \in \{1, \dots, 2N+1\}^m} \Lambda_{\mu,z}(z) P_{\nu,z}(\zeta^{-1} \circ z) D_\nu^z(\zeta) U_{1,z}.$$

We observe that

$$(4.23) \quad |V_z(\zeta) - V_{\bar{z}}(\zeta)| \leq cd_{\bar{z}}(\bar{z}, z)^\alpha, \quad \forall z, \bar{z} \in E_0, \zeta \in \Omega_0,$$

for some constant c . Indeed, (4.23) follows from (3.8), the fact the function $z \mapsto P_{\nu,z}$ is of class C_B^α by Lemma 3.5, and the estimate

$$(4.24) \quad |D_\nu^z P_z^k u(\zeta) - D_\nu^{\bar{z}} P_{\bar{z}}^k u(\zeta)| \leq cd_{\bar{z}}(\bar{z}, z)^\alpha.$$

At this point, the proof of (4.20) is analogous to the one of (4.17). Indeed it suffices to observe that

$$D_\sigma(z)I_{1,z} - D_\sigma(\bar{z})I_{1,\bar{z}} = A_1(z, \bar{z}) + A_2(z, \bar{z}) + A_3(z, \bar{z}),$$

where

$$\begin{aligned}
 A_1(z, \bar{z}) &= \int_{\Omega_0 \setminus \Omega_M(\bar{z}, z)} (\Gamma_z(z, \zeta) V_z(\zeta) - \Gamma_{\bar{z}}(\bar{z}, \zeta) V_{\bar{z}}(\zeta)) d\zeta, \\
 A_2(z, \bar{z}) &= \int_{\Omega_M(\bar{z}, z)} (\Gamma_z(z, \zeta) - \Gamma_{\bar{z}}(\bar{z}, \zeta)) V_{\bar{z}}(\zeta) d\zeta, \\
 A_3(z, \bar{z}) &= \int_{\Omega_M(\bar{z}, z)} \Gamma_z(z, \zeta) (V_z(\zeta) - V_{\bar{z}}(\zeta)) d\zeta,
 \end{aligned}$$

and to conclude as before, by using Proposition 3.1. Thus, (4.15) is proved. We omit the proof of (4.16) since it is essentially analogous.

The case $k = 2$. Let $b, f \in C_b^\alpha(\Omega)$, with $0 < \alpha < 1$. Since, by assumption (1.10), $b \in C^1(\Omega)$, we have

$$L_B u(z) - L_{\bar{z}} u(z) = (b - P_{\bar{z}}^1 b) \partial_{y_1} u(z) = O(d_{\bar{z}}(\bar{z}, z)), \quad \text{as } z \rightarrow \bar{z}.$$

In this case, a much simpler choice of the frozen operators provides an approximation of L_b of the same order. Indeed, for convenience, let us denote by $z = (x, y, t)$ a point of Ω . For fixed $\bar{z} \in \Omega$, we define

$$L^{(\bar{z})} = \Delta_x + (b(\bar{z}) + x_1 - \bar{x}_1) \partial_y - \partial_t.$$

Then, $L^{(\bar{z})}$ is a Hörmander operator which, in the case $N = 1$, up to a straightforward change of variables, coincides with the Kolmogorov operator in (1.5). Moreover, we have

$$\begin{aligned} L_b u(z) - L^{(\bar{z})} u(z) &= (b(z) - b(\bar{z}) - (x_1 - \bar{x}_1)) \partial_y u(z) \\ &= O(d^{(\bar{z})}(\bar{z}, z)), \quad \text{as } z \rightarrow \bar{z}, \end{aligned}$$

where $d^{(\bar{z})}$ denotes the control distance associated to $L^{(\bar{z})}$.

Given a cut-off function φ , we represent the solution u in terms of the fundamental solution $\Gamma^{(\bar{z})}$ of $L^{(\bar{z})}$:

$$(u\varphi)(z) = \int_{\Omega} \Gamma^{(\bar{z})}(z, \zeta) L^{(\bar{z})}(u\varphi)(\zeta) d\zeta.$$

Since $L^{(\bar{z})}(u\varphi) \in C_0(\Omega)$, it is standard to prove that $u \in C_b^{1, \bar{\alpha}}(\Omega)$, for every $\bar{\alpha} \in]0, 1[$. Moreover, it is not difficult to adapt the previous arguments and to show that $u \in C_b^{2, \bar{\alpha}}(\Omega)$, for every $\bar{\alpha} \in]0, \alpha[$. $\square$

Remark 4.3. Theorem 1.3 holds true if we assume that u is a locally Lipschitz continuous, strong solution to (1.1) instead of a classical solution. We recall that u is a strong solution to (1.1) if it has weak derivatives and equation (1.1) is satisfied almost everywhere. In order to justify our claim, it suffices to remark that the proof of Theorem 1.3 is based only on the representation formula (4.2) and on the boundedness of the first-order derivatives of the solution.

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